

The turbulent wake

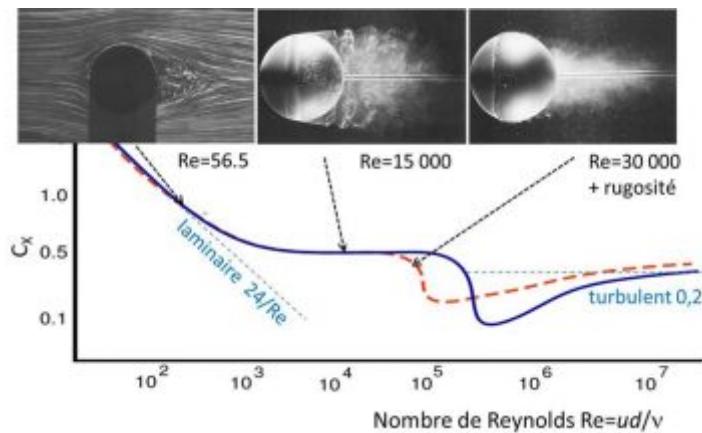


Figure 1. Removal of a wake behind a sphere at different Reynolds numbers, represented on the drag coefficient curve: at $Re=56.5$, the flow detaches downstream but remains laminar; at $Re=15,000$, the flow is turbulent but its boundary layer remains laminar; at $Re=30,000$, the boundary layer is destabilized by a thin wire stuck on the sphere, thus reproducing the behaviour at $Re \approx 106$ too fast to be visualized. The rod visible on the right is the support of the sphere [Source: Documents Werlé Onera and van Dyke, Album of fluid motion]

The Figure opposite illustrates, on the emblematic case of a sphere, how the wake changes from a laminar regime to an increasingly turbulent regime as the Reynolds number increases.

For a moderate value, $Re=56.5$, we see that the trajectories are detached from the sphere but only on its downstream part, unlike the symmetrical flow obtained at very low Re (see Figure 2 of the article). This detachment still has little effect on the drag coefficient, whose law remains close to that of the laminar regime in $24/Re$, as can be seen on the blue curve.

At the larger Re , the flow consists upstream of a thin sheared layer enveloping the surface of the sphere, called the **boundary layer**, while downstream it is turbulent. The wake, visualized by the injection of smoke, occupies a large cross section, corresponding to a fairly large drag coefficient $Cx \approx 0.5$, a quasi-constant value when Re is between 10^3 and 10^5 .

Above $Re=105$, the drag drops strangely and then stabilizes at a value close to 0.2. This lower Re drop can be reproduced by disturbing the boundary layer with a roughness or an annular wire stuck on the sphere, which is the case of the photo on the right, at $Re=30000$. We then see that the boundary layer itself becomes turbulent, and this has the effect of delaying its disbonding, leading to a narrower wake. This is the explanation for the drop in the drag coefficient.

It is to produce such a reduction in drag that golf balls are not smooth but honeycombed: for a speed of 50 m/s and a diameter of 4 cm, we have $Re=2 \cdot 10^5$, and we can see on the figure that for this value, the drag for a rough sphere (red dotted line) is indeed lower than that of a smooth sphere.

The case of very large Reynolds numbers can be seen as the borderline case of a very low viscosity. A fluid with strictly zero viscosity would slide frictionlessly over the obstacle, and its flow would represent a kind of ideal where the upstream and downstream would be symmetrical and laminar, and where the drag would cancel each other out. However, this theoretical limit is never reached, which is the paradox of D'Alembert. The explanation lies in the fact that the viscosity, however low it may be, always produces a boundary layer, thicker downstream than upstream, which imposes an asymmetry between upstream and downstream. In addition, at very low viscosity, the boundary layer becomes unstable and feeds wake turbulence, leading to the non-zero drag boundary value observed at very large Reynolds numbers.

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