

# Planck's theory

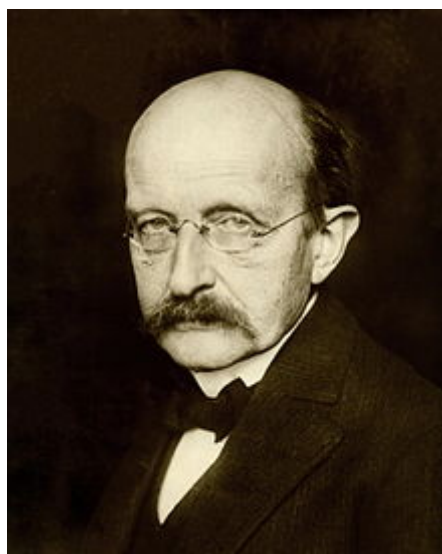


Figure 1. Physicist Max Planck.

All attempts to explain the behaviour of the spectral density  $\rho_\nu(\nu, T)$  (shown in Figure 2 of the article) with the concepts of classical physics have failed. The theory is most easily derived in the particular case of an empty cubic cavity, with perfectly reflective walls, of unit volume. A phenomenon of standing waves similar to that of a vibrating rope fixed at its two ends then occurs. In a cavity, a **mode** characterizes a particular stationary electromagnetic wave given by its frequency  $\nu$ , its propagation speed (or celerity)  $c$ , its propagation direction, and by the direction and amplitude of its electric field. The possible frequencies are discrete and are referred to as normal frequencies. The latter form a dense set of modes characterized by a density (number of modes per unit volume of the cavity)  $D(\nu)$  proportional to  $\nu^2$ .

In classical physics, the average energy  $\langle E(\nu) \rangle$  of a standing wave comparable to a harmonic oscillator [1] is simply  $k_B T$  where  $k_B$  is the Boltzmann constant and  $T$  the absolute temperature. This expresses the fact that at thermal equilibrium the energy is distributed evenly in all accessible modes. The energy spectral density  $\rho_\nu(\nu, T)$ , given by the product  $D(\nu) \langle E(\nu) \rangle$ , is then proportional to  $k_B T \nu^2$ . This result of classical physics, represented by the dotted line in Figure 2, is called Rayleigh's model. By integrating over all frequencies from 0 to  $\nu$ , it leads to a total energy density  $\rho$  proportional to  $k_B T \nu^3$ . It grows without limit at high frequency  $\nu$  leading to an infinite total energy, therefore unacceptable. This the so-called "ultraviolet catastrophe" led Max Planck to formulate his quantum hypothesis mentioned above.

The energy of a mode is quantized: the only possible energy values of a frequency mode are  $E_n = n h \nu$ , where  $n$  is a positive or zero integer and  $h$  the Planck constant, which is identical for all oscillators. An oscillator in an energy state  $E_n$  is equivalent to  $n$  quanta or **photons of energy  $h \nu$** . The energy exchanges between the wave and matter (atoms, molecules) are thus carried out in a discrete way by exchanging photons. By quantifying the energy of each oscillator in this way, a statistical calculation [2] [3]

$$\langle E(\nu) \rangle = \frac{h \nu}{\exp\left(\frac{h \nu}{k_B T}\right) - 1} = \langle N \rangle h \nu$$

shows that the average energy of each oscillator is now given by

where  $\langle N \rangle$  is the average number of photons of the mode with energy  $h \nu$  at temperature  $T$ . For low-frequency modes ( $h \nu$  small in front of  $k_B T$ ), this expression reduces to the classical law  $\langle E \rangle = k_B T$ . These modes contain many photons. However the energy of the modes drops exponentially as  $h \nu$  clearly exceeds  $k_B T$ , and  $\langle N \rangle$  becomes much smaller than 1. By multiplying the mean energy of the modes by their density  $D(\nu)$ , we obtain the famous **Planck formula** :

$$\rho_v(\nu, T) = \frac{8\pi h}{c^3} \frac{\nu^3}{\exp\left(\frac{h\nu}{k_B T}\right) - 1}$$

of frequency  $\nu$  has been shown in Figure 2 for various temperatures. This formula is in remarkable agreement with experiences. It also allows to experimentally determine the two constants  $h/c^3$  and  $h/k_B$  from which the value of the Planck constant  $h = 6.62617 \times 10^{-34}$  Js and the Boltzmann constant  $k_B = 1.38066 \times 10^{-23}$  J/K is derived. This leads to an unexpected evaluation of the Avogadro number  $N_A = R/k_B = 6.0221 \times 10^{23} \text{ mol}^{-1}$ , in excellent agreement with its value obtained by other methods ( $R$  is the constant of perfect gases, read [pressure, temperature, heat](#)).

By integrating on all frequencies we obtain the energy density of the black body  $\rho = \alpha T^4$  as well as the radiated  $\Phi$  flux per unit area of the black body given by Stefan-Boltzmann's law  $\Phi = \sigma T^4$ , with  $\sigma = 2\pi^5 k_B^4 / (15c^2 h^3)$ . This leads to results that are remarkably in line with experience. Note that if we express the spectral energy density as a function of wavelength, Planck's

$$\rho_\lambda(\lambda, T) = \frac{8\pi hc}{\lambda^5} \frac{1}{\exp\left(\frac{hc}{\lambda k_B T}\right) - 1}$$

theory leads to [\[4\]](#)

The curves in Figures 4 and 5 are representative of this

function and Wien's law results from the position of its maximum.

This Planck's theory published in 1901 was the first step towards the quantification of light, further confirmed by Einstein in 1905 who explicitly introduced the notion of photon to explain the photoelectric effect. The existence of energy levels in atoms and molecules was then proposed by Niels Bohr in 1913 to explain the line spectra. These results were later incorporated in the wider Quantum Mechanics theory which governs all microscopic phenomena. This formalism was gradually developed in the 1920s, with contributions from Louis De Broglie, Erwin Schrödinger and Werner Heisenberg.

## References and notes

[\[1\]](#) Each mode is characterized by a solution of the type  $E_z = E_0 \cos(\omega t) \sin k_x x$ , where  $z$  is the direction of the electric field  $\mathbf{E}$ ,  $\omega = 2\pi\nu$  is the pulsation of the wave and  $k_x = 2\pi/\lambda$  is the wave vector along the propagation direction  $x$ . A solution of this type is similar to that of a mechanical oscillator: a mass  $m$  suspended on a spring whose restoring force is proportional to the elongation  $x$  and whose movement is described by a harmonic oscillation of the form  $x = x_0 \cos \omega t$ .

[\[2\]](#) Elie Belorizky, *Initiation to Quantum Mechanics*, Dunod, Paris, 2003.

[\[3\]](#) Elie Belorizky, *Introduction to Statistical Mechanics*, Presse Universitaire de Grenoble, Grenoble sciences collection, 1992.

[\[4\]](#) We have the correspondence  $\nu = c/\lambda$  and the equality of the energy densities :  $\rho_\lambda(\lambda, T) |d\lambda| = \rho_\nu(\nu, T) |d\nu|$ , which leads to  $\rho_\lambda(\lambda, T) = \rho_\nu(\nu, T) |d\nu/d\lambda| = \rho_\nu(c/\lambda, T) c/\lambda^2$ .

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